

$\mathcal{F}_S = \frac{\mathcal{F}}{\mathcal{F}_0}$ multilocal functionals

\mathcal{F}_S functionals that vanish on E_S

Observation: Take a vector field $x \in \Gamma(T\mathcal{E})$, which is also multilocal

$\langle dL, x \rangle$ vanishes on E_S (by definition)

Assumption: Assume that all elements of \mathcal{F}_0 are this way

↳ Henneaux: proof of that assumption fulfilled in examples of physical interest

Marc Henneaux, Lectures on the antifield-BRST formalism for gauge theories.
Nuclear Physics B - Proceedings Supplements, 18(1):47–105, December 1990.

Let V_{loc} denote local vector fields
 V multilocal
 V_{reg} regular } all compactly supported

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$-d_{dL}(V) \subset \mathcal{F}_0$ by definition

By assumption: $d_{dL}(V) \subseteq \mathcal{F}_0$



Extend δ to ΛV by using the graded Leibniz rule

$$\text{assumption} \quad \delta^2 = 0$$

$$\text{compute } H^0 = \mathcal{F}/\delta(V) \stackrel{!}{=} \mathcal{F}/\mathcal{F}_0 = \mathcal{F}_S \quad \text{functionals on the solution space}$$

Compute H^{-1} : $\text{Ker } \delta$ consists of v. fields such that

$$\delta(x) = -\langle dL, x \rangle \equiv 0$$

" $\partial_x L(f)$, $f=1$ on supp x

x is a symmetry!

$\text{Im } \delta: X \wedge Y \in \Lambda^2 V$

$$\delta(X \wedge Y) = -\langle dL, X \rangle Y + \langle dL, Y \rangle X$$

such symmetry vanishes on the solution space E_S , hence is called trivial

$\Rightarrow H^{-1}$ contains non-trivial local symmetries

Examples: • scalar field has no non-trivial local symmetries

• Yang-Mills theories do have symmetries

Sufficient condition for H^{-1} to be trivial: P_Q is normally hyperbolic (e.g. d'Alembertian)

P_Q denotes the diff. op. induced by $L^{(2)}(Q)$ linearized EOM, $Q \in \mathcal{E}$

↳ Proof: Fredenhagen, KR:

Fredenhagen, K., Rejzner, K.: Batalin-Vilkovisky formalism in the functional approach to classical field theory. Commun. Math. Phys. 314(1), 93–127 (2012).

$P_Q: \mathcal{E} \rightarrow \mathcal{E}^* \subset \mathcal{E}'$ normally hyperbolic $\Rightarrow \exists$ retarded and advanced green functions:

$$\Delta_Q^R / \Delta_Q^A: \mathcal{E}' \rightarrow \mathcal{E} \text{ s.t. } \Delta_Q^R \circ P_Q|_{\mathcal{E}} = \text{id}_{\mathcal{E}}$$

$\Delta_Q^R(f)$ ↗ image in \mathcal{E}'^*

$$P_Q \circ \Delta_Q^A = \text{id}_{\mathcal{E}'^*}$$

$\Delta_Q^A(f)$ ↙ support properties

Define: $\Delta_q \doteq \Delta_q^R - \Delta_q^A : \mathcal{E}_c^* \rightarrow \mathcal{E}$ Pauli-Jordan function
(commutator function)

\hookrightarrow induces an integral kernel $\Pi_q \in \Gamma_c' (\mathcal{E}^* \otimes \mathbb{M}^2 \rightarrow \mathbb{M}^2)$

$q \mapsto \Pi_q$ is a generalized bivector

This is in fact a Poisson bracket. Non-trivial thing to show: Jacobi identity ($(\Delta_q^A)^T = \Delta_q^A$ and that $L^{(2)}$ is symmetric)

Jakobs "Eichbrücken in der klassischen Feldtheorie" (Appendix B)

<http://www-library.desy.de/preparch/desy/thesis/desy-thesis-09-009.pdf> (AQFT Hamburg webpage)

Define: Poisson bracket (Peierls) : $\langle F, G \rangle_q \doteq \langle \Pi_q, dF(q) \otimes dG(q) \rangle$

$F, G \in \mathcal{E}$

Peierls, R.E.: The commutation laws of relativistic field theory. Proc. R. Soc. Lond. Ser. A. Math. Phys. Sci. **214**(1117), 143–157 (1952)

Observation: Extends to $\Lambda \mathcal{U}$, but the space of multilocal functionals is not closed under L, \cdot, \cdot . Requires completion.

For example: use microcausal functionals / v. fields

\hookrightarrow conditions on singularity structure of derivatives $\Lambda \mathcal{U}_{mc}$

For a review see e.g.:

Kasia Rejzner. Perturbative Algebraic Quantum Field Theory: An Introduction for Mathematicians. Springer, March 2016.

Classical alg model: $\mathcal{O} \mapsto (\Lambda \mathcal{U}_{mc}(\mathcal{O}), L, \cdot, \cdot, \delta)$ contains local non-linear observables

Also possible:

$\mathcal{O} \mapsto (\Lambda \mathcal{U}_{reg}(\mathcal{O}), L, \cdot, \cdot, \delta)$ well-defined but boring, only contains linear local observables and their products

More graded geometry:

Observation: $\Lambda \mathcal{U} = \Gamma(T^* \mathcal{E} \sqcup \mathcal{E})$ shifted cotangent bundle

Observation: $\Lambda \mathcal{U}$ as the space of multivector fields comes with the Schouten bracket:

- $\{X, FY\} := -\partial_X F$, $F \in \mathcal{E}, X \in \mathcal{U}$
- $\{X, Y\} := -[X, Y]$, $X, Y \in \mathcal{U}$
- graded Leibniz rule

Aka "antibracket"

Notation: For a vector field $X \in \mathcal{U}$, denote: $X = \int X(x) \frac{\delta}{\delta \varphi^*(x)}$

$$\{X, YY\} = \left\langle \frac{\delta X}{\delta \varphi}, \frac{\delta Y}{\delta \varphi^*} \right\rangle - \left\langle \frac{\delta X}{\delta \varphi^*}, \frac{\delta Y}{\delta \varphi} \right\rangle$$

\hookrightarrow right derivative $\left. \frac{\delta}{\delta \varphi}\right|_{\varphi^*}$ antifield

$$\{X, FY\} = - \int X(x) \frac{\delta F}{\delta \varphi(x)} \equiv - \int X(x) \frac{\delta F}{\delta \varphi^*(x)} = - \partial_X F$$